

Super Yang-Mills Theory in 10+2 dims. Another Step Toward M-theory

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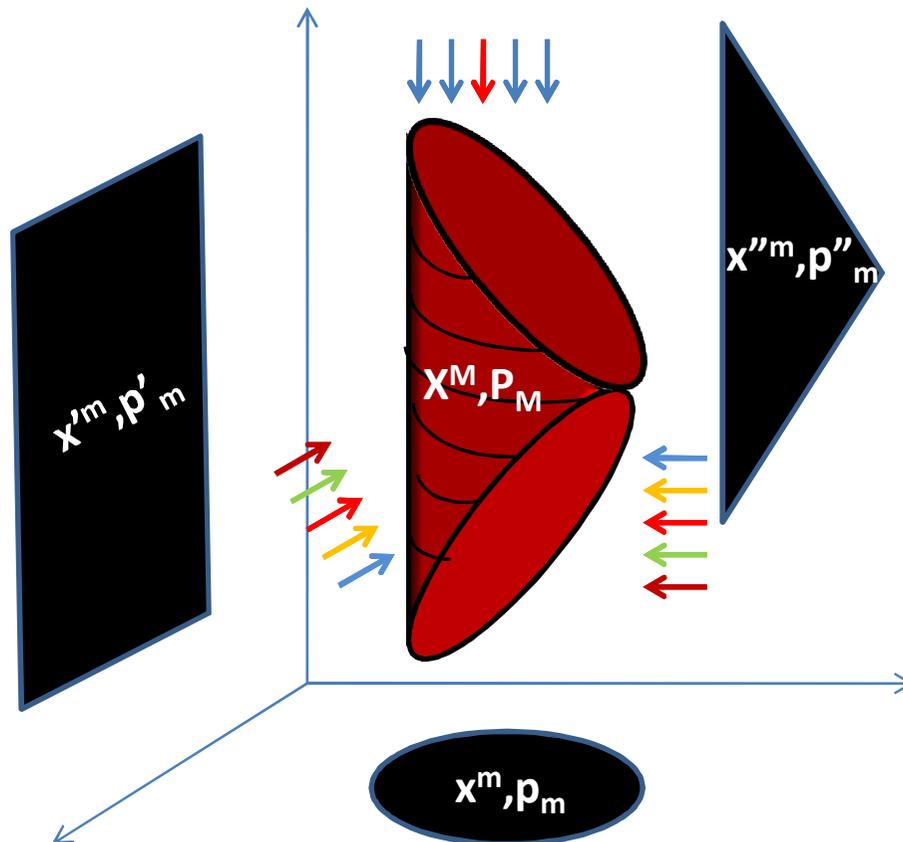
http://physics.usc.edu/~bars/homepage/moscow2009_bars.pdf

- SYM exists only in 2+1, 3+1, 5+1 and 9+1 dimensions. I will report² on a new path which enlarges this horizon. I will show that the new theory is the mother of the N=4 SYM in 3+1 dims, the N=1 SYM in 9+1 dims, and M(atrrix) theory, and others. The new theory is developed in the context of 2T-physics.
- Sakharov, who was one of the first to entertain the notion of two times, would have enjoyed what I now call 2T-physics.
- Strong hints for 2T-physics came from M-theory (IB -1995):
Extended SUSY of M-theory is really a SUSY in 12 dimensions
 $\{Q_{32}, Q_{32}\} = Z_{[2]} + Z_{[6]+}$, Q_{32} real Weyl spinor \rightarrow **(10+2) signature!**
But if this implies 2 times, how does one remove the ghosts?
- 2T-physics developed by finding the fundamental solution to this ghost problem, and related causality problem. The answer is a gauge symmetry in phase space X^M, P_M . Phase space gauge symmetry is reminiscent of U-duality in M-theory (electric-magnetic).
- After a crash review on 2T-physics, I will explain the new SYM theory.

2T-Physics as a unifying framework for 1T-physics

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- 2T-physics is a ghost-free general framework that correctly describes all physics.
- 2T-physics and usual 1T-physics are related, but 2T-physics unifies a larger set of phenomena that 1T-physics is unable to predict, but is only able to verify.



The relation between 2T-physics and 1T-physics can be described by an analogy : Object in the room (4+2 dim. phase space, X^M, P_M) and its shadows on walls (3+1 dim **many phase spaces, x^m, p_m**).

Observers like us are stuck on the “walls” (3+1 dims.), no privilege to be in the room (4+2). We interpret the shadows as different dynamical systems (1T formalism).

One (2T) to many (1T's). Predict many relations among the shadows (dualities, symmetries). This is systematically missed information in 1T-physics approach.

1) 1T-physics is incomplete !!!

2) Is 2T-physics more suitable for fundamentals?

2T-physics principles in a nutshell⁴

Basic principle: Position-Momentum symmetry at every instant, for all motion for all physics (?)

Sp(2,R) gauge symmetry, local on worldline $X^M(\tau)$, $P_M(\tau)$

3 generators: $Q_{11}(X,P)$, $Q_{22}(X,P)$, $Q_{12}(X,P)=Q_{21}(X,P)$

Generalize worldline action for $x^\mu(\tau)$, $p^\mu(\tau)$

$$\partial_\tau x^\mu p_\mu - \frac{1}{2} e p_\mu p_\nu \eta^{\mu\nu}$$

Example: spinless particle $A^{ij}(\tau)$ is Sp(2,R) gauge potential

$$\mathcal{L}_{2T} = \partial_\tau X^M P_M - \frac{1}{2} A^{ij} Q_{ij}(X, P)$$

simplest example: $Q_{ij}(X,P) = (X \cdot X, P \cdot P, X \cdot P)$
first class constraints $Q_{ij}(X,P) = 0$: requires Sp(2,R) singlets

Generalize?
 nontrivial soln.
 and no ghosts :
ONLY Sp(2,R) !!
Only 2T !!

Physical sector: only gauge invariant motion is allowed (**shadows**)

Nontrivial solutions exist only with 2 times! No less and no more!

The “shadows” are in 1 less space and 1 less time: $[(d-1)+1]$ (gauge fixed)

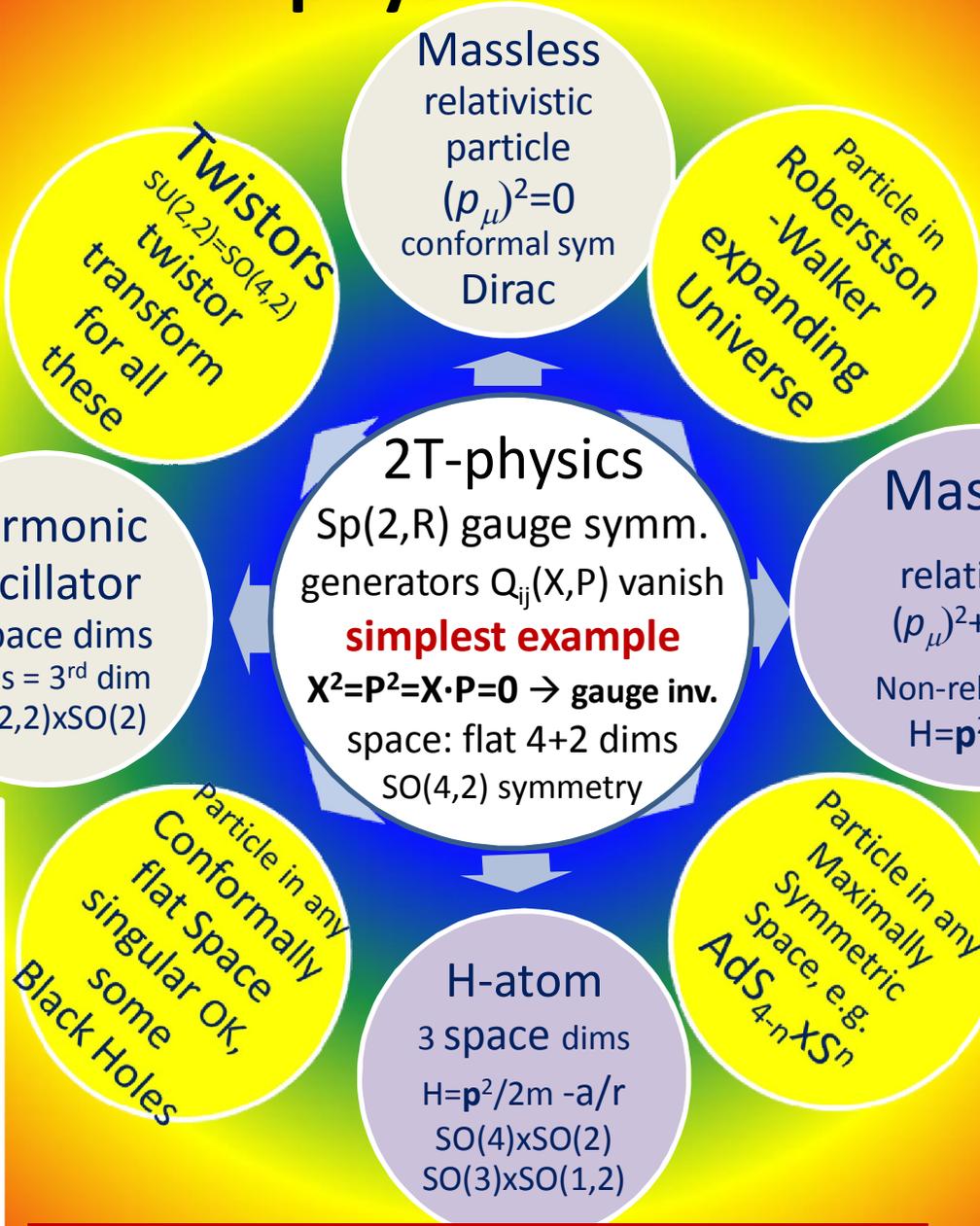
In the simple example, spacetime η_{MN} : flat $d+2$ dims., $SO(d,2)$ global symmetry

Shadows from 2T-physics → hidden info in 1T-physics ⁵

Hidden Symm.
 $SO(d,2)$, ($d=4$)
 $C_2=1-d^2/4 = -3$
 singleton

Emergent
 spacetimes
 and emergent
 parameters:
 mass,
 couplings,
 curvature, etc.

• Holography:
 These emergent
 holographic
 shadows are only
 some examples of
**much broader
 phenomena.**



Free or interacting
 systems, with or
 without mass, in
 flat or curved
 3+1 spacetime
**Analogy: shadows
 on walls**

2T-physics predicts
hidden symmetries
 and dualities (with
 parameters) among
 the "shadows".
**1T-physics misses
 these phenomena.**

These emerge in 2T-field theory as well

Rules for 2T field theory, spins=0, 1/2, 1

Impose Sp(2,R) singlet condition !!

Use BRST approach for Sp(2,R). Like string field theory: I.B.+Kuo, hep-th/0605267
I.B. hep-th/060645

Flat space

$$S_{kin} = \int d^{d+2} X \delta(X^2) \left[\frac{1}{2} \bar{\Phi} D^2 \Phi + \frac{i}{2} \bar{\Psi} X \not{D} \Psi + h.c. - \frac{1}{4} F_{MN} F^{MN} \Omega^{\frac{2(d-1)}{d-2}} \right], \quad \Omega \text{ is dilaton}$$

There is explicit X^M , no translation invariance, only **SO(d,2) spacetime invariance**. This SO(d,2) becomes conformal symmetry in the “conformal shadow”, but a hidden SO(d,2) symmetry in other shadows.

$$S_{yukawa} = \int d^{d+2} X \delta(X^2) \Omega^{-\frac{d-1}{d-2}} \left[y (\bar{\Psi}_L X \Psi_R) \Phi + h.c. \right], \quad \Psi_{L,R} \text{ spinors of } SO(d,2)$$

Double the size spinor as SO(d-1,1) + Fermionic gauge sym.

$$S_{scalars} = \int d^{d+2} X \delta(X^2) V(\Omega, \Phi), \quad V(\Omega, H) = \Omega^{\frac{2d}{d-2}} V\left(1, \frac{\Phi}{\Omega}\right)$$

Homogeneous $V(\Omega, \Phi)$
Only dimensionless couplings among scalars

$$S_{anomalies} \sim \int d^{d+2} X \delta(X^2) \varepsilon^{M_1 M_2 M_3 \dots M_{d+2}} (X_{M_1} \hat{c}_{M_2} \ln \Omega) (A_{M_3 \dots M_{d+2}})$$

$$\delta S(\Phi) = 2\gamma \int d^{d+2} X \delta\Phi \left\{ \begin{array}{l} \delta(X^2) [\partial^2 \Phi - V'(\Phi)] \\ + 2\delta'(X^2) [X \cdot \partial \Phi + \frac{d-2}{2} \Phi] \end{array} \right\}$$

dynamical eq. $\mathbf{p^2 + \dots = 0}$
kinematic eqs. $\mathbf{X^2=0, X.P+P.X=0}$

Minimizing the action gives two equations, so get **all 3 Sp(2,R) constraints for each field, including interaction !!**

New gauge symmetries + kinematic equations (\Leftrightarrow Sp(2,R)), eliminate all ghosts!!

Gravity in 2T-physics Field Theory

Gauge symmetry and consistency with $Sp(2, R)$ lead to a unique action in $d+2$ dims, with **no parameters at all**.

$$a = \frac{(d-2)}{8(d-1)}$$

$$S^0 = \gamma \int d^{d+2} X \sqrt{G} \left\{ \begin{array}{l} \delta(W) \left[\Omega^2 R(G) + \frac{1}{2a} \partial \Omega \cdot \partial \Omega - V(\Omega) \right] \\ + \delta'(W) \left[\Omega^2 (4 - \nabla^2 W) + \partial W \cdot \partial \Omega^2 \right] \end{array} \right\}$$

Pure gravity has triplet of fields:
 $G_{MN}(X)$, metric
 $\Omega(X)$, dilaton
 $W(X)$, replaces X^2

It has unique coupling to matter: scalars, spinors & vectors.

Imposes **severe constraints on scalar fields** coupled to gravity.

$$S_{shadow}(g, \phi, s_i) = \int d^d x \sqrt{-g} \left(\begin{array}{l} \frac{1}{2a} g^{\mu\nu} \partial_\mu \phi \partial_\nu \phi - \frac{1}{2} g^{\mu\nu} \partial_\mu s_i \partial_\nu s_i \\ + (\phi^2 - a s_i^2) R - V(\phi, s_i) \end{array} \right)$$

Local scale symm $\lambda(x)$ comes from general coordinate symm in $d+2$. Can choose dilaton $\phi(x)$ arbitrarily, e.g. a constant

$$\phi_0^2 = \frac{1}{16\pi G_d}$$

=> **Gravitational scale.**

$$g'_{\mu\nu} = e^{2\lambda} g_{\mu\nu}, \quad \phi' = e^{-\frac{d-2}{2}\lambda} \phi, \quad s'_i = e^{-\frac{d-2}{2}\lambda} s_i$$

Prediction from 2T-physics: The gravitational constant is determined by the vacuum values of all scalar fields. It increases after every cosmic phase transition at the scales of inflation, GUT, SUSY, electroweak. **Effect on cosmology !!**

Super Yang-Mills in 10+2 dimensions

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General SUSY Field Theory, for N=1,2,4, in 4+2 dimensions
done: IB + Y-C.Kuo hep-th/ 0702089, 0703002, 0808.0537

Usual N=4 SYM in d=4 is the conformal shadow from 4+2

12D theory

$$S = 2\gamma \int (d^{10+2}X) \sqrt{G} \delta(W(X)) L(A_m^a(X), \lambda_a^a(X))$$

vector 10+2, spinor 32 (Weyl)

Note G,W,Ω general gravity background

$$L(A, \lambda)_{W, \Omega} = -\frac{1}{4} \Omega^{3/2} F_{mn}^a F_{m'n'}^a G^{mm'} G^{nn'} + \frac{i}{2} \left[\bar{\lambda}^a V \bar{D} \lambda^a + \bar{\lambda}^a \bar{D} V \lambda^a \right]$$

$$V \equiv \gamma^m V_m, \quad \bar{D} \equiv \bar{\gamma}^m D_m = \bar{\gamma}^m \left(\partial_m + \frac{1}{4} \omega_m^{ij} \gamma_{ij} + A_m^a t^a \right)$$

forms of V_m, G_{mn} follow from $\text{Sp}(2, R)$

non-dynamical
background W, Ω

$$V_m = \frac{1}{2} \partial_m W, \quad V \cdot V = W$$

$$G_{mn} = \frac{1}{2} \nabla_m \nabla_n W, \quad (G^{mn} V_m \partial_n + 4) \Omega = 0$$

Homothety: Lie derivative $\mathcal{L}_V G^{MN} = -2G^{MN}$

SUSY is possible due to special
gamma matrix identity in 10+2 dims

$$(\gamma^{ik})_{(\alpha\beta} \left(\gamma_i^j \right)_{\gamma)\delta} + (\gamma^{ij})_{(\alpha\beta} \left(\gamma_i^k \right)_{\gamma)\delta} = \frac{1}{6} \eta^{kj} (\gamma^{il})_{(\alpha\beta} (\gamma_{il})_{\gamma)\delta}$$

$$\begin{aligned} & 2f_{abc} (V_n \bar{\epsilon} \gamma^{qn} \lambda^a) (V^p \bar{\lambda}^b \gamma_{qp} \lambda^c) \delta(W) \\ &= \frac{1}{6} f_{abc} (\bar{\lambda}^b \gamma_{il} \lambda^a) (\bar{\epsilon} \gamma^{il} \lambda^c) W \delta(W) = 0 \end{aligned}$$

similar identity also in (3+2), (4+2), (6+2)

conserved current $\partial_m(\bar{\epsilon}J^m) = 0 \Leftrightarrow$ SUSY $\delta_\epsilon A_m, \delta_\epsilon \lambda_\alpha$
 $\epsilon_\alpha = 32$ spinor

$$\bar{\epsilon}^\alpha J_\alpha^m = \delta(W) \sqrt{G} \Phi^{3/4} F_{pq}^a V_n (\bar{\epsilon} \gamma^{pqn} \bar{\gamma}^m \lambda)$$

SUSY possible only iff ϵ_α satisfies SUSY condition

$$\left[-(\bar{\epsilon} \gamma^m \bar{\gamma}^{pqn} \lambda) V_n \partial_m \ln \Phi^{\frac{d-4}{d-2}} + V_n (D_m \bar{\epsilon}) \gamma^{pqn} \bar{\gamma}^m \lambda \right]_{W=0}$$

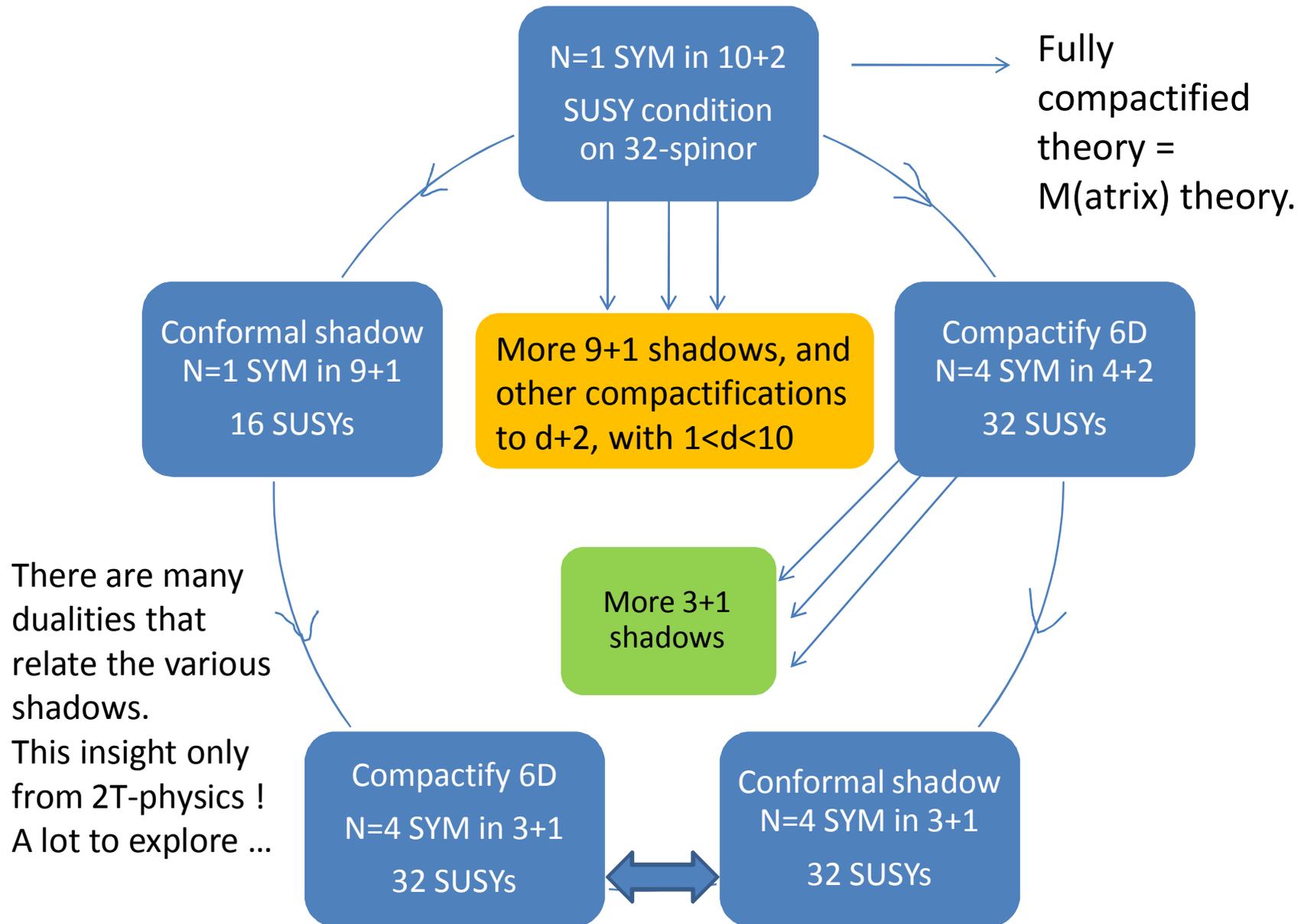
$$= (V^p U^q - V^q U^p), \quad \text{any } U^q(X)$$

$$D_m \epsilon_\alpha = 0 \quad \text{and} \quad \bar{\epsilon} \gamma^m \partial_m \ln \Phi^{\frac{d-4}{d-2}} = 0.$$

For most background geometries such $\epsilon(X)$ can be found with only 16 independent components.

But there are special cases with 32 components. For example, dimensionally reduce 10+2 to (4+2)+(6+0), then we obtain 32 component ϵ which corresponds to N=4 SYM in 4+2 dimensions, which in turn has N=4 SYM in 3+1 dimensions in the conformal shadow.

10+2 SYM as parent of N=4 SYM in 3+1, and a web of dualities



Status of 2T-physics

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- Local $Sp(2,R)$ → 2T-physics, a principle in CM & QM:
Seems to work generally to produce 1T Hamiltonians for particle dynamics, including spin, supersymmetry, backgrounds of all types, including gravity, E&M, etc.. **A new unification of 1T systems into classes that belong to the same 2T system**, and brings out hidden symmetries related to extra dims.
- Field Theory, The Standard Model & Gravity in 4+2 dimensions,
In the “conformal shadow” in 3+1 dims. agree structurally with usual SM and GR, but include some new constraints that provide new phenomenological guidance for physics at the LHC and in Cosmology (e.g. $-\frac{1}{12} s^2 R$ is required!!)
- Beyond the Standard Model
GUTS, SUSY, higher dims; all have been elevated to 2T-physics in d+2 dimensions.
Strings, branes; tensionless, and twistor superstring, 2T OK. Tensionful incomplete.
M-theory; expect 11+2 dimensions → $Osp(1|64)$ global SUSY, S-theory.
- New non-perturbative technical tools – a lot more to do here !!
Emergent spacetimes and dynamics; unification; holography; duality; hidden sym. Expect to be useful for non-perturbative analysis of field theory, including QCD. (analogs of AdS-CFT, others ...). Path integral approach for quantum field theory directly in d+2 dimensions will be useful. (still under development).

IB+Chen+Quelin, 0705.2834 0802.1947,

Do you need 2T? YES!

Hidden information in 1T-physics is revealed by 2T-physics (shadows)

1T-physics on its own is not equipped to capture these hidden symmetries and dualities, which actually **exist**.

1T-physics needs the additional guidance, so 1T-physics is definitely incomplete.



"Of course the elements are earth, water, fire and air. But what about chromium? Surely you can't ignore chromium."

2T-physics seems to be a promising idea on a new direction of higher dimensional **unification**.

extra 1+1 are LARGE, also not Kaluza-Klein, **not hidden**.

Different shadows are different perspectives, so you can "see" extra dims. indirectly by proper interpretation.

A lot more remains to be done with 2T-physics. Predictions at every scale of physics are expected from hidden dualities and symmetries by using the more powerful tools in future research ...

**2T-physics works
in the known world so far**

**... and through work in progress
we hope to the extend
its domain of validity
to solve the remaining mysteries!!**

The End